

「名古屋大学HPC計算科学連携研究プロジェクト」
シンポジウム 5/10/2011

地球流体乱流の数値解析

木村 芳文

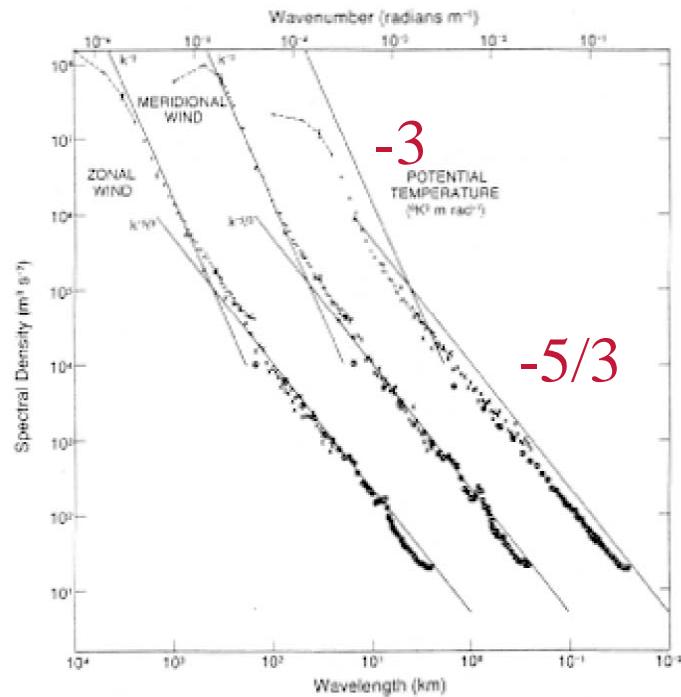
名古屋大学多元数理科学研究所

collaboration:

Jackson R. Herring

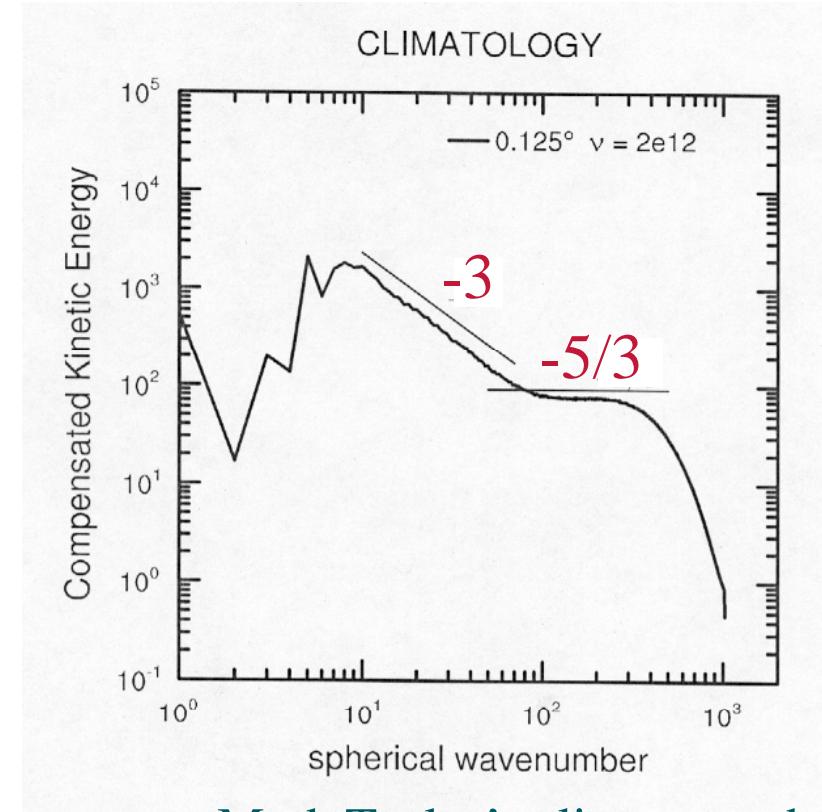
National Center for Atmospheric Research

Transition in Energy Spectrum for Rotating and Stratified turbulence



Nastrom-Gage's atmospheric observation (1985)
(JAS **42** 950-960.)

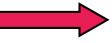
Both stratification and rotation are essential



Mark Taylor's climate model simulation (2008)
(CCSM project at NCAR)

- 3 : enstrophy cascade for Quasi-Geostrophic turbulence (~2D)
- 5/3 : Kolmogorov turbulence (3D)

Transition in Energy Spectrum for Stratified turbulence

Observations: $k_z^{-2} \sim -3$  $k^{-5/3}$
(in the ocean) Garret-Munk spectrum Kolmogorov spectrum

Munk (1981), Garrett *et.al* (1981)
transition wavenumber: $k_c \sim \sqrt{N^3/\varepsilon}$ (Ozmidov scale)

Theory: Munk (1981), Garrett *et.al* (1981), Lumley (1964), Holloway (1983)
All support the Ozmidov scale for transition

Simulation: Carnevale, Briscoline & Orlandi (2001) LES at 128^3
Yoshida, Ishihara & Kaneda (2002) LES up to 512^3
Waite & Bartello (2004) DNS + hyperviscosity
~ Ozmidov for transition

(Waite & Bartello (2004) for the review)

Navier–Stokes equation with the Boussinesq approximation

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u} + \theta \hat{\mathbf{z}} + \mathbf{F}$$

$$\frac{\partial \theta}{\partial t} + (\mathbf{u} \cdot \nabla) \theta = \kappa \nabla^2 \theta - N^2 w$$

$$\nabla \cdot \mathbf{u} = 0$$

where

$\mathbf{u} = (u, v, w)$: velocity

θ : temperature fluctuations

$N^2 = \frac{g \alpha}{T_0} \frac{\partial \bar{T}}{\partial z}$: Brunt - Väisälä frequency

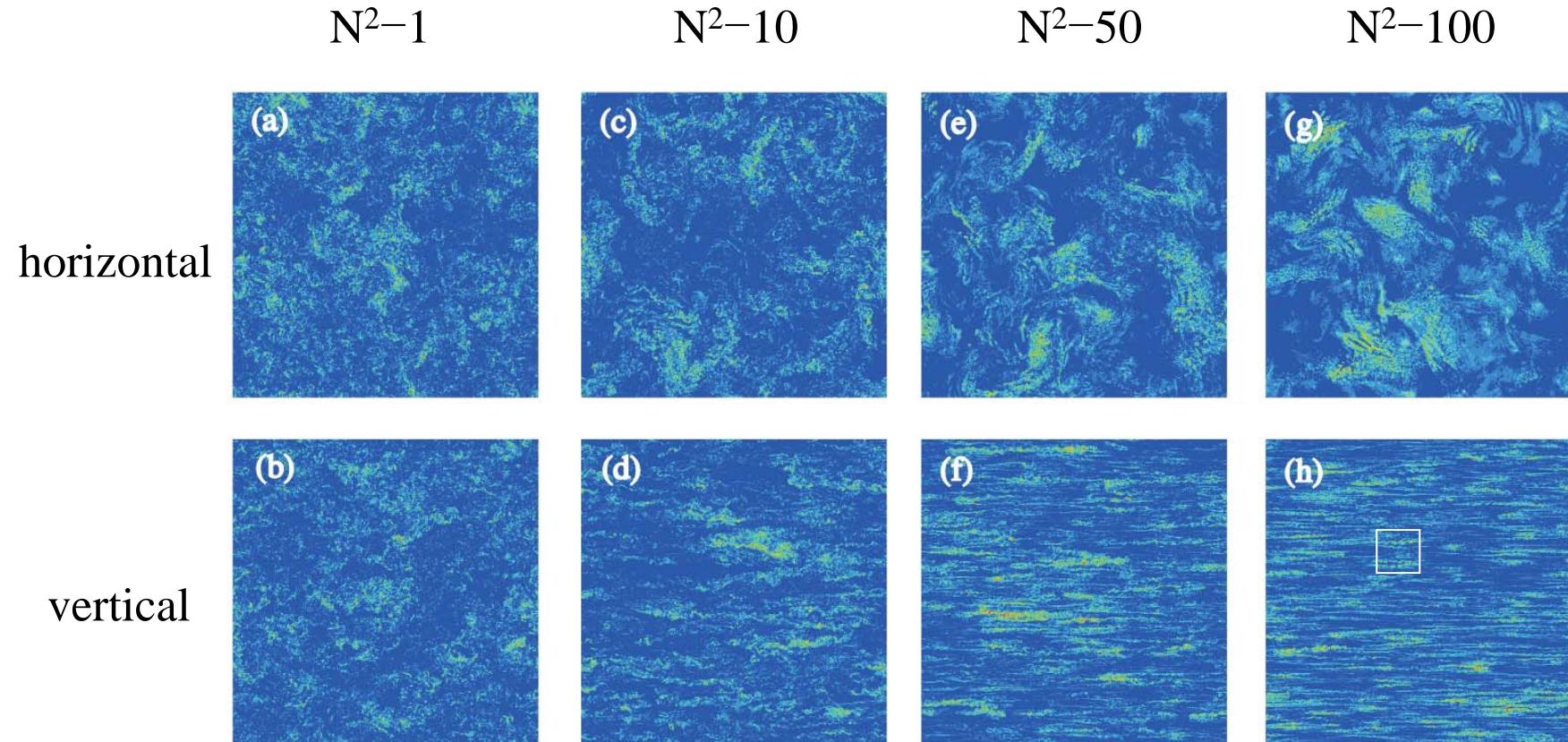
\mathbf{F} : Forcing (horizontal)

Numerical Methods

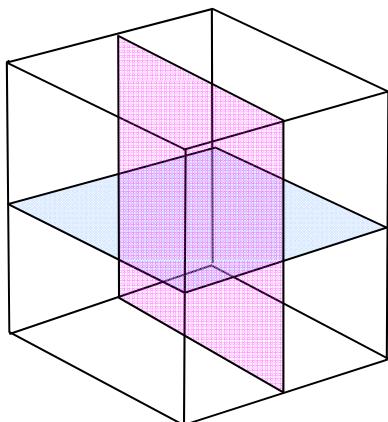
- ◆ forced simulations
- ◆ 2π -periodic box with 1024^3 grid points ($R_\lambda \sim 300$)
- ◆ 3rd order time-marching scheme
- ◆ Initial energy spectrum : $E(k) = 0$
- ◆ Force horizontal velocity components
- ◆ Add red noise to modes within a wave number band
 $(k_f \sim 5)$

Solving Ornstein-Uhlenbeck process (2nd order stochastic ODEs)

Enstrophy contours

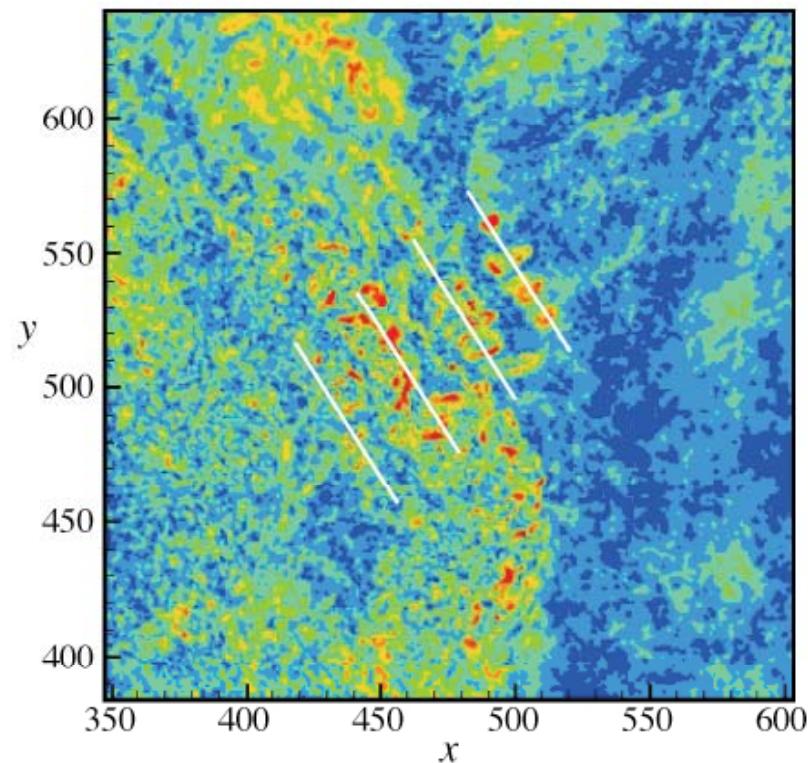
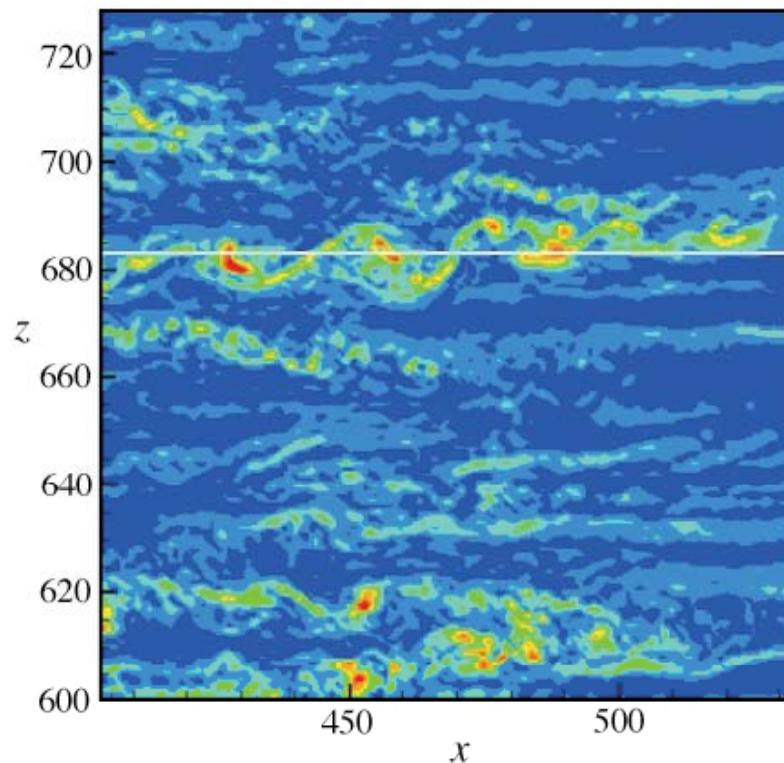


As N^2 becomes large;



- ◆ large scale clusters and elongated streaks appear in the horizontal plane.
- ◆ thin layers and wedge structures develop in the vertical plane.

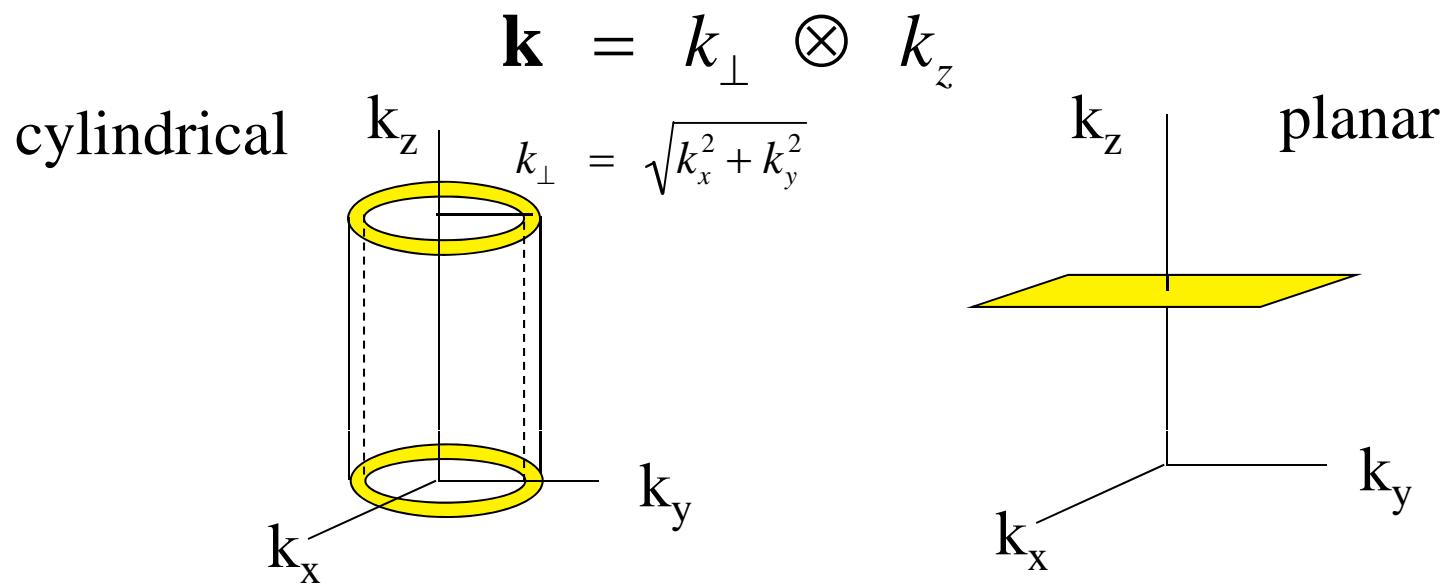
Enstrophy contours (blow-up)



- ◆ Kelvin-Helmholz billows are observed in the vertical.
- ◆ The billows are not single rollers and chopped in the horizontal.

Characteristics of stratified turbulence

- ◆ Composite of “waves” and “turbulence”
 - “*Craya-Herring decomposition*” to separate waves and turbulence
- ◆ Highly anisotropic
 - Need suitable averaging



“Craya-Herring” decomposition

$$\nabla \cdot \mathbf{u} = 0$$

incompressibility



$$\mathbf{k} \cdot \tilde{\mathbf{u}} = 0$$

$\tilde{\mathbf{u}}$ is spanned by two independent vectors perpendicular to \mathbf{k}

$$\begin{aligned}\mathbf{e}_1(\mathbf{k}) &= \frac{\mathbf{k} \times \hat{\mathbf{z}}}{\|\mathbf{k} \times \hat{\mathbf{z}}\|} = \frac{1}{\sqrt{k_x^2 + k_y^2}} \begin{pmatrix} k_y \\ -k_x \\ 0 \end{pmatrix} \\ \mathbf{e}_2(\mathbf{k}) &= \frac{\mathbf{k} \times \mathbf{k} \times \hat{\mathbf{z}}}{\|\mathbf{k} \times \mathbf{k} \times \hat{\mathbf{z}}\|} = \frac{1}{\sqrt{k_x^2 + k_y^2 + k_z^2} \sqrt{k_x^2 + k_y^2}} \begin{pmatrix} k_z k_x \\ k_z k_y \\ -(k_x^2 + k_y^2) \end{pmatrix} \\ \mathbf{e}_3(\mathbf{k}) &= \frac{\mathbf{k}}{\|\mathbf{k}\|} = \frac{1}{\sqrt{k_x^2 + k_y^2 + k_z^2}} \begin{pmatrix} k_x \\ k_y \\ k_z \end{pmatrix}\end{aligned}$$

orthnormal coordinates

$$\tilde{\mathbf{u}}(\mathbf{k}) = \phi_1 \mathbf{e}_1(\mathbf{k}) + \phi_2 \mathbf{e}_2(\mathbf{k})$$

$$\phi_1 = \tilde{\mathbf{u}}(\mathbf{k}) \cdot \mathbf{e}_1(\mathbf{k})$$

$$\begin{aligned}&= \frac{1}{\sqrt{k_x^2 + k_y^2}} (k_y \tilde{u} - k_x \tilde{v}) \\ &= \frac{i}{\sqrt{k_x^2 + k_y^2}} \tilde{\omega}\end{aligned}$$

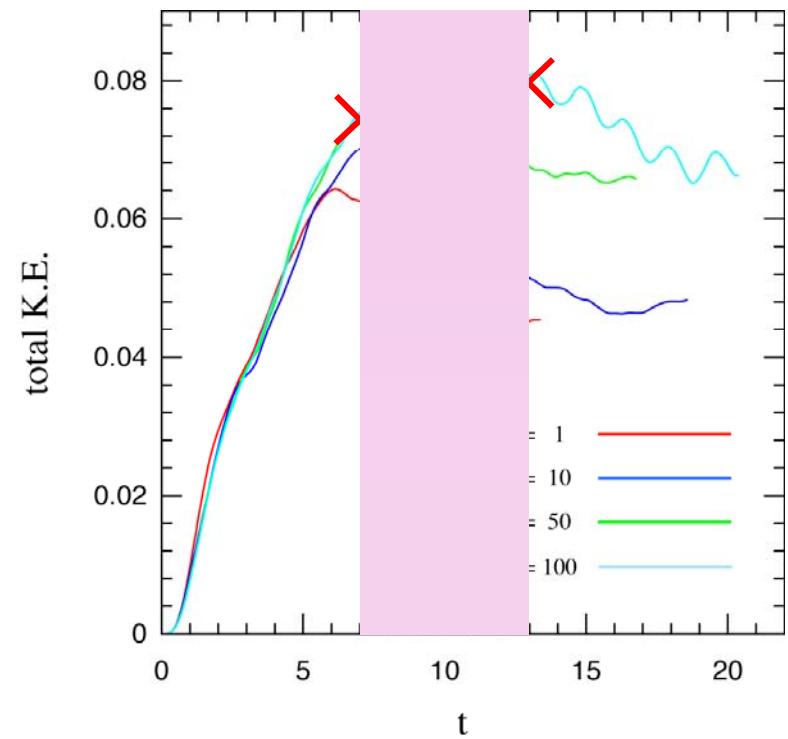
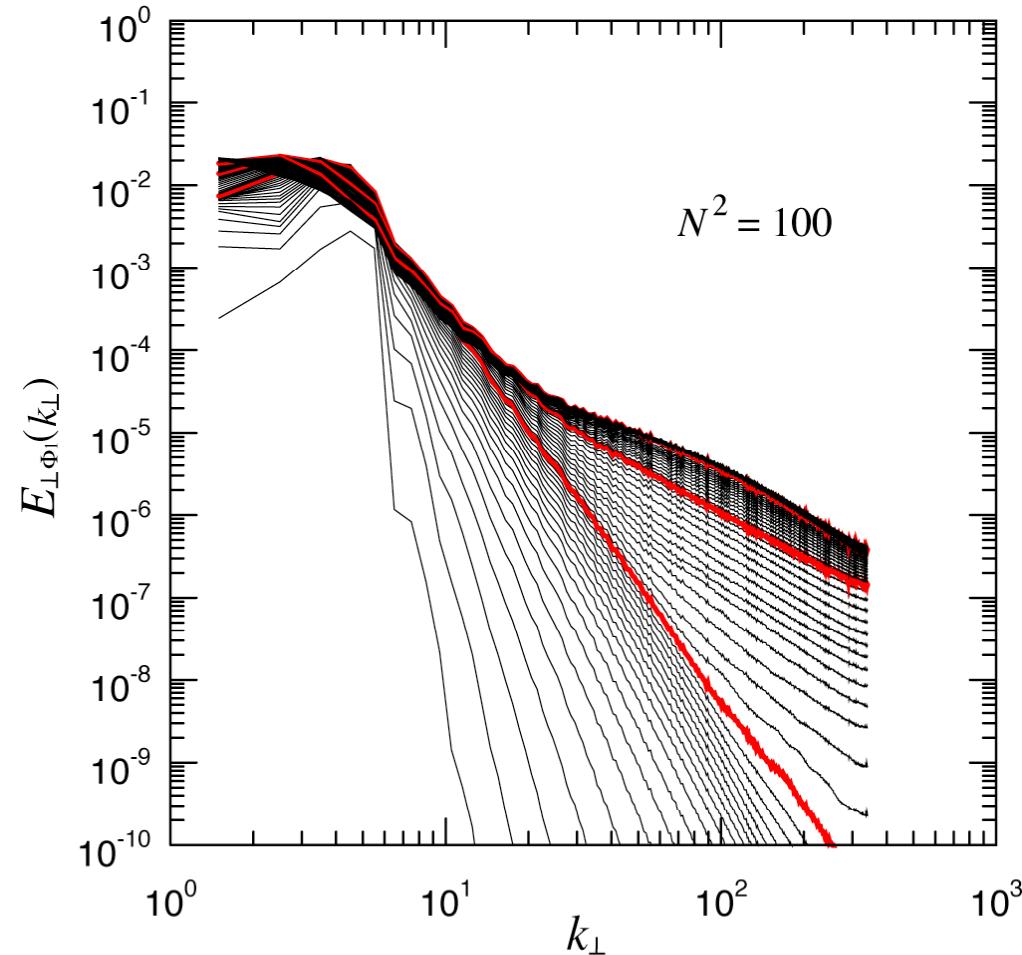
(vortex, rotation)

$$\phi_2 = \tilde{\mathbf{u}}(\mathbf{k}) \cdot \mathbf{e}_2(\mathbf{k})$$

$$= \frac{\sqrt{k_x^2 + k_y^2 + k_z^2}}{\sqrt{k_x^2 + k_y^2}} \tilde{w}$$

(wave, divergence)

History of Φ_1 energy spectra ($N^2=100$)



First, steep spectrum ($\sim k^{-3}$) develops then small scales rise.

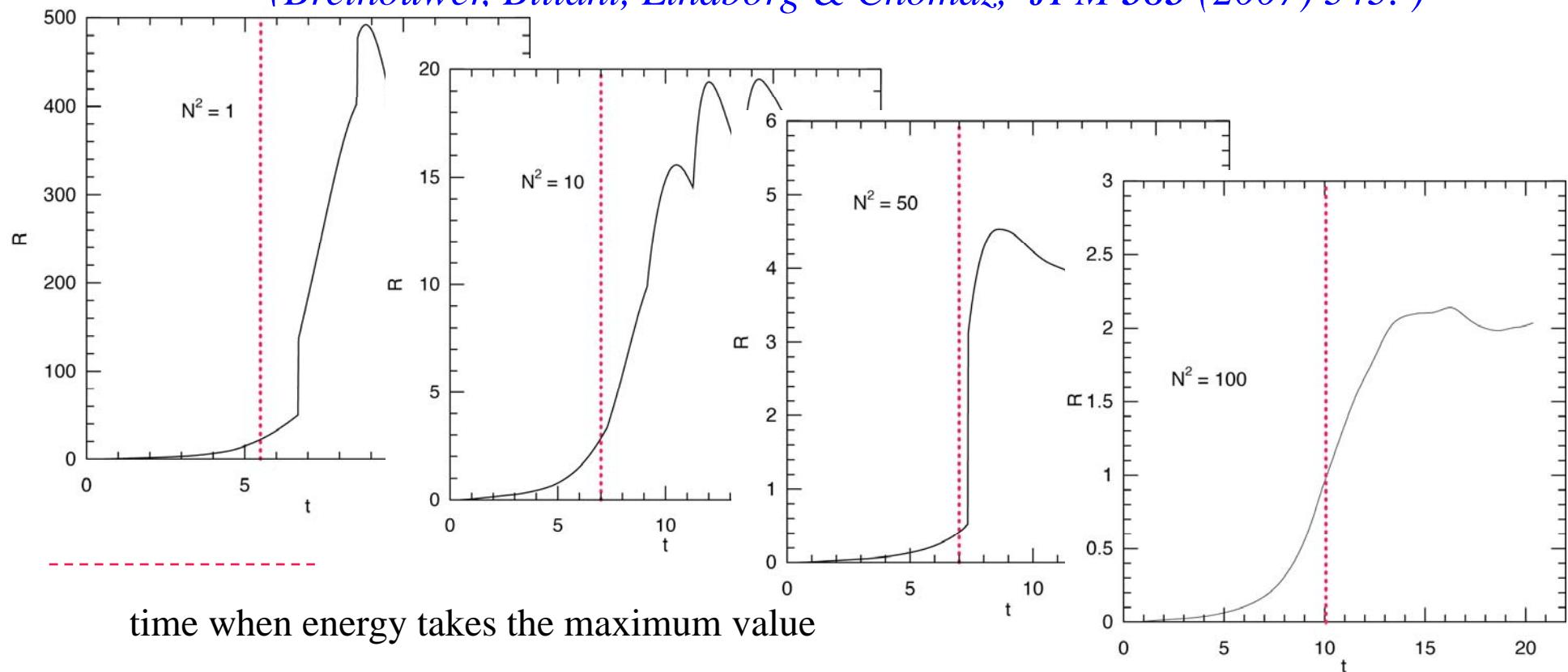
History of buoyancy Reynolds number

$$R = \text{Fr}_h^2 \text{Re} = \frac{\varepsilon}{\nu N^2} \rightarrow \left[\sqrt{\frac{\varepsilon}{\nu N^2}} \middle/ \left(\frac{\nu^3}{\varepsilon} \right)^{1/4} \right]^{4/3} = [L_O/L_K]^{4/3}$$

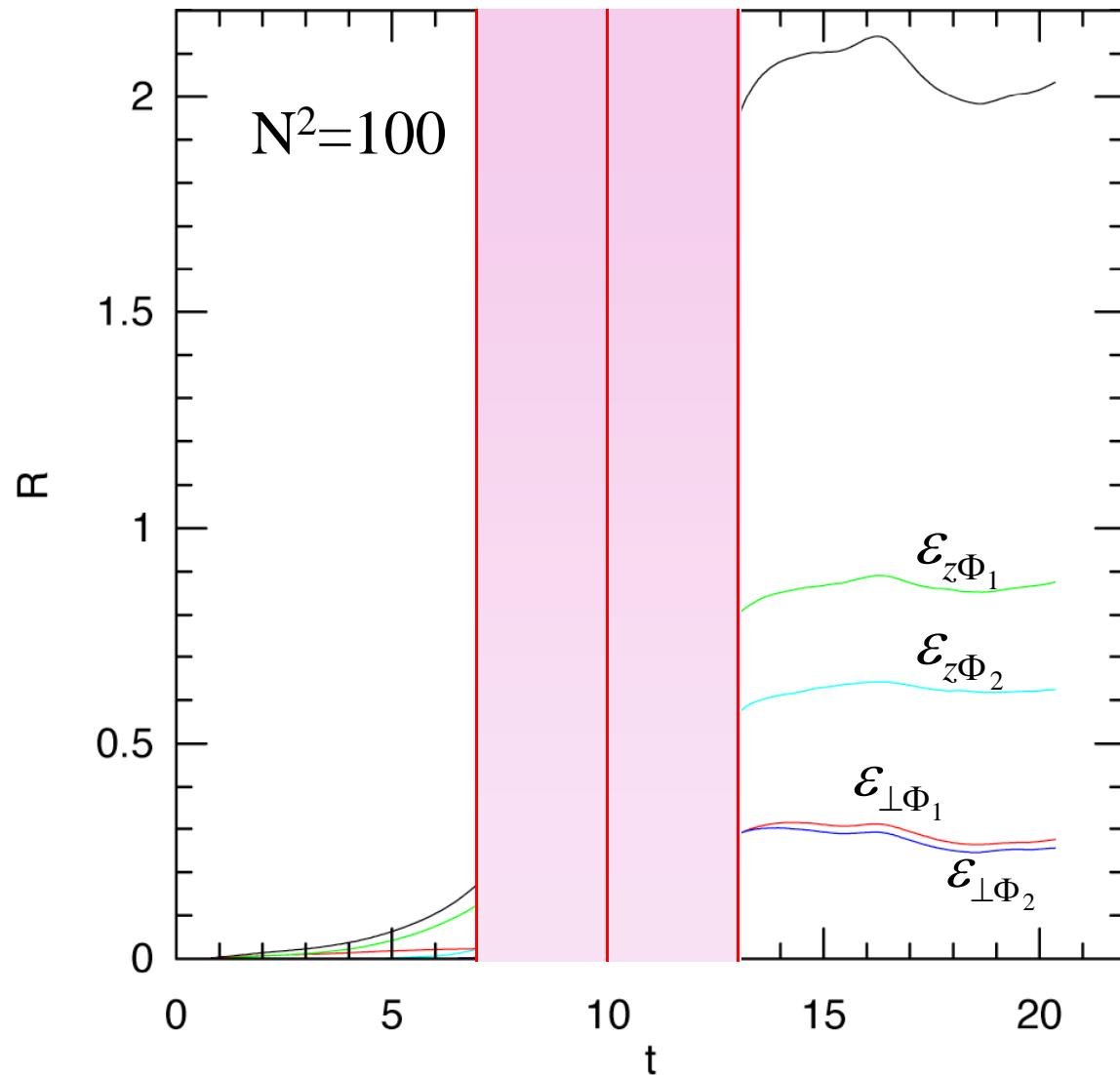
L_O : Ozmidov scale L_K : Kolmogorov scale

$R < 1$: steep spectrum, $R > 1$: -5/3.

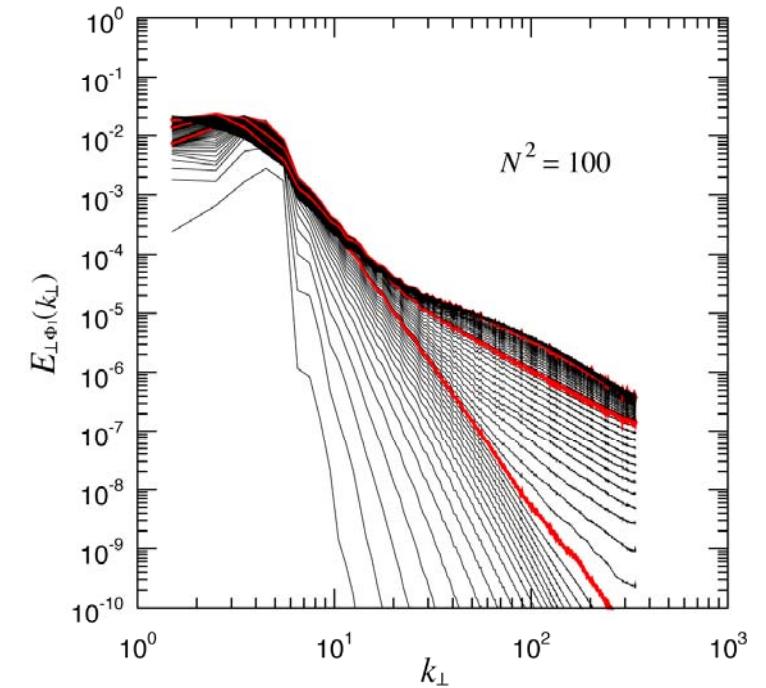
(Brethouwer, Billant, Lindborg & Chomaz, JFM 585 (2007) 343.)



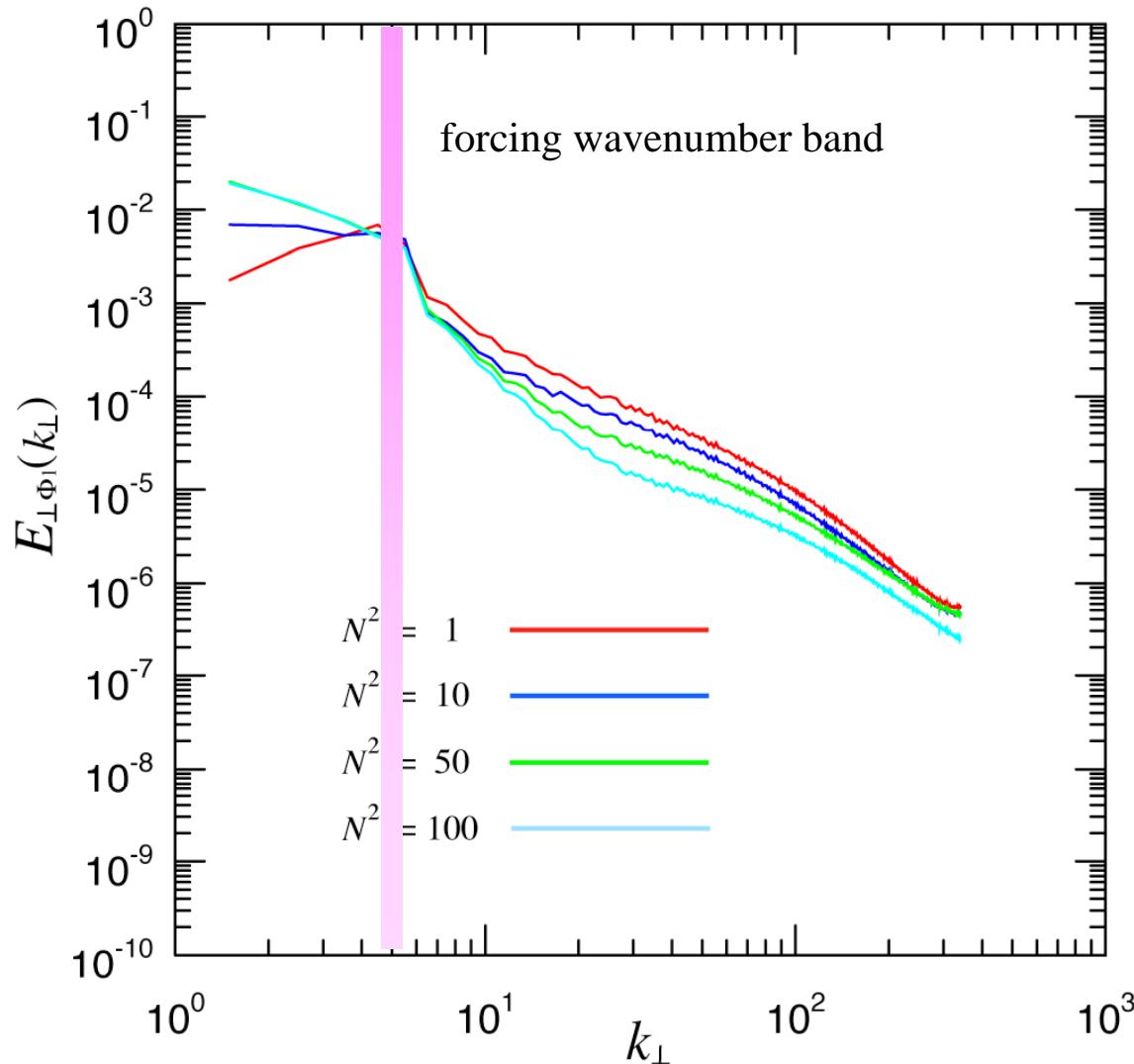
History of buoyancy Reynolds number



$$R = \text{Fr}_h^2 \text{Re} = \frac{\varepsilon}{\nu N^2}$$



$\Phi_1(\underline{k})$ spcetra for various N



- ◆ -5/3 for small scales
(different coefficient)
- ◆ -3 for large scales
(same coefficient)
- ◆ There's a sharp transition

How to scale them ?

guidelines:

- ◆ The Kolmogorov const. is universal. (Sreenivasan 85)
- ◆ The coefficient for large scales doesn't depend on N

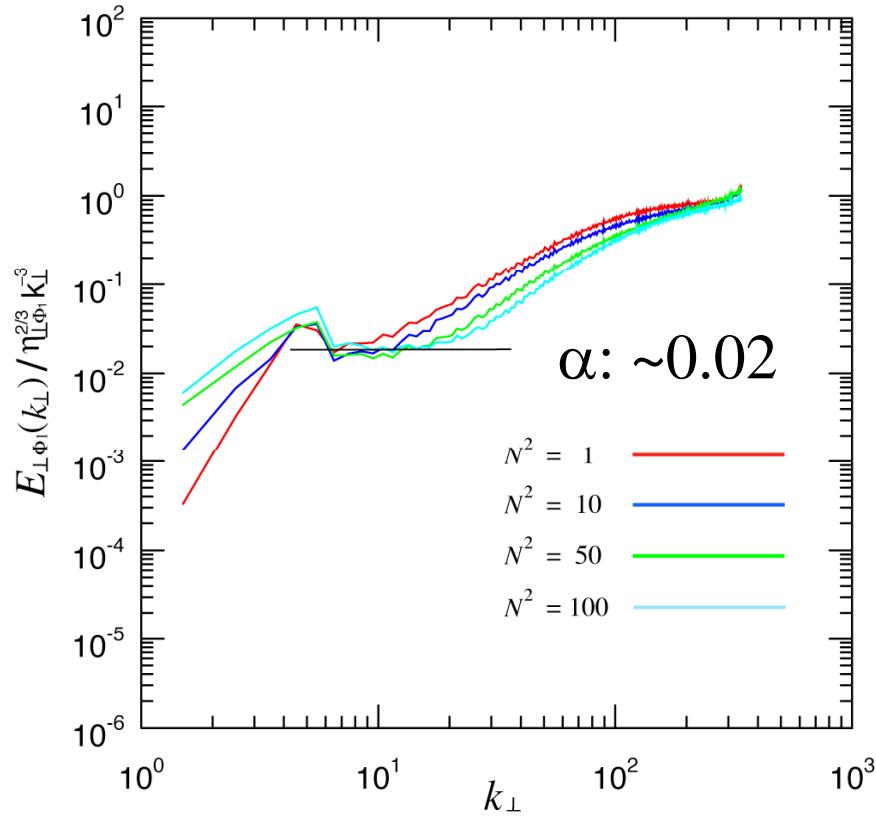
Anisotropy in dissipation

$$\begin{aligned}
 \varepsilon &= 2\nu \int_0^\infty \mathbf{k}^2 E(\mathbf{k}) d\mathbf{k} = 2\nu \int_0^\infty (k_\perp^2 + k_z^2) (E_{\Phi_1} + E_{\Phi_2}) d\mathbf{k} \\
 &= \underbrace{\varepsilon_{\perp\Phi_1} + \varepsilon_{\perp\Phi_2}}_{\text{horizontal dissipation}} + \underbrace{\varepsilon_{z\Phi_1} + \varepsilon_{z\Phi_2}}_{\text{vertical dissipation}}
 \end{aligned}$$

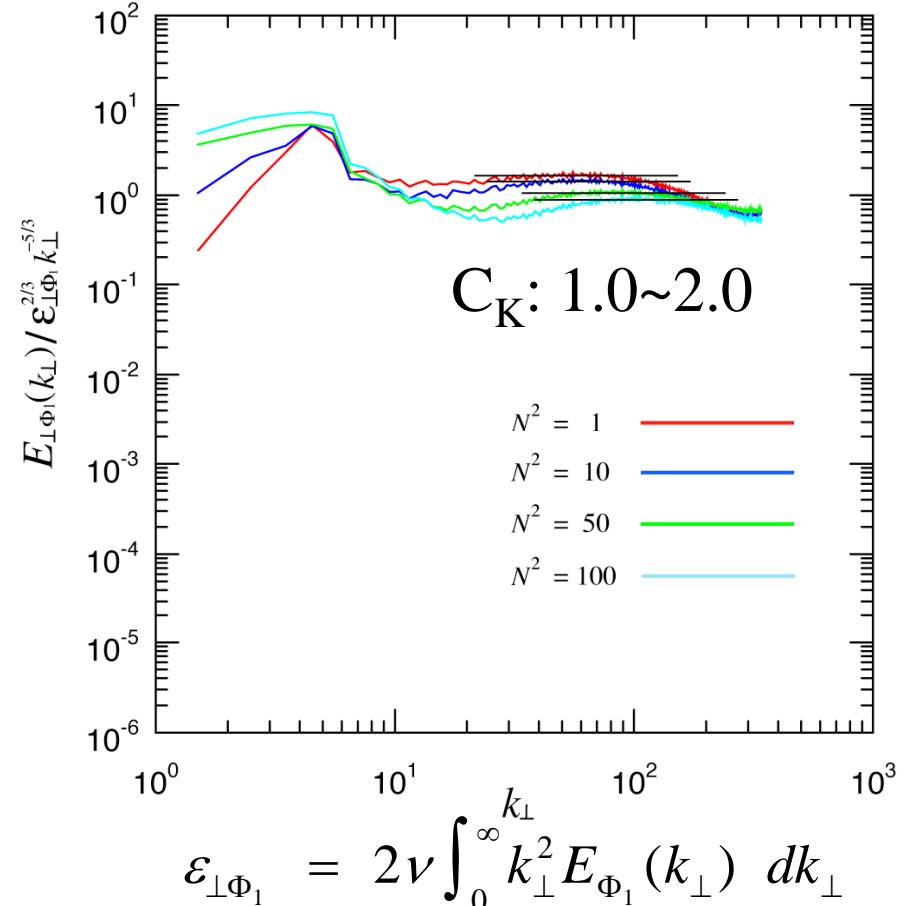
$\langle |\partial u / \partial z|^2 \rangle \longrightarrow \frac{N^2}{\langle |\partial u / \partial z|^2 \rangle}$
 (local Richardson number)

N^2	$\varepsilon_{\perp\Phi_1}$	$\varepsilon_{\perp\Phi_2}$	$\varepsilon_{z\Phi_1}$	$\varepsilon_{z\Phi_2}$	ε
1	1.63×10^{-3}	1.72×10^{-3}	9.28×10^{-4}	9.45×10^{-4}	5.23×10^{-3}
10	1.24×10^{-3}	1.46×10^{-3}	1.12×10^{-3}	1.18×10^{-3}	5.01×10^{-3}
50	1.09×10^{-3}	1.10×10^{-3}	2.07×10^{-3}	1.74×10^{-3}	6.00×10^{-3}
100	6.61×10^{-4}	5.93×10^{-4}	2.55×10^{-3}	1.68×10^{-3}	6.04×10^{-3}

Compensated spectra



$$\eta_{\perp\Phi_1} = 2\nu \int_0^\infty k_\perp^4 E_{\Phi_1}(k_\perp) dk_\perp$$



Conjecture of $\Phi_1(k_\perp)$ spectra

$$E_{\Phi_1}(k_\perp) = \begin{cases} \alpha \eta_{\perp\Phi_1}^{2/3} k_\perp^{-3} & (k_\perp < k_c) \\ C_K \varepsilon_{\perp\Phi_1}^{2/3} k_\perp^{-5/3} & (k_\perp > k_c) \end{cases}$$

where

$$\eta_{\perp\Phi_1} = 2\nu \int_0^\infty k_\perp^4 E_{\Phi_1}(k_\perp) dk_\perp$$

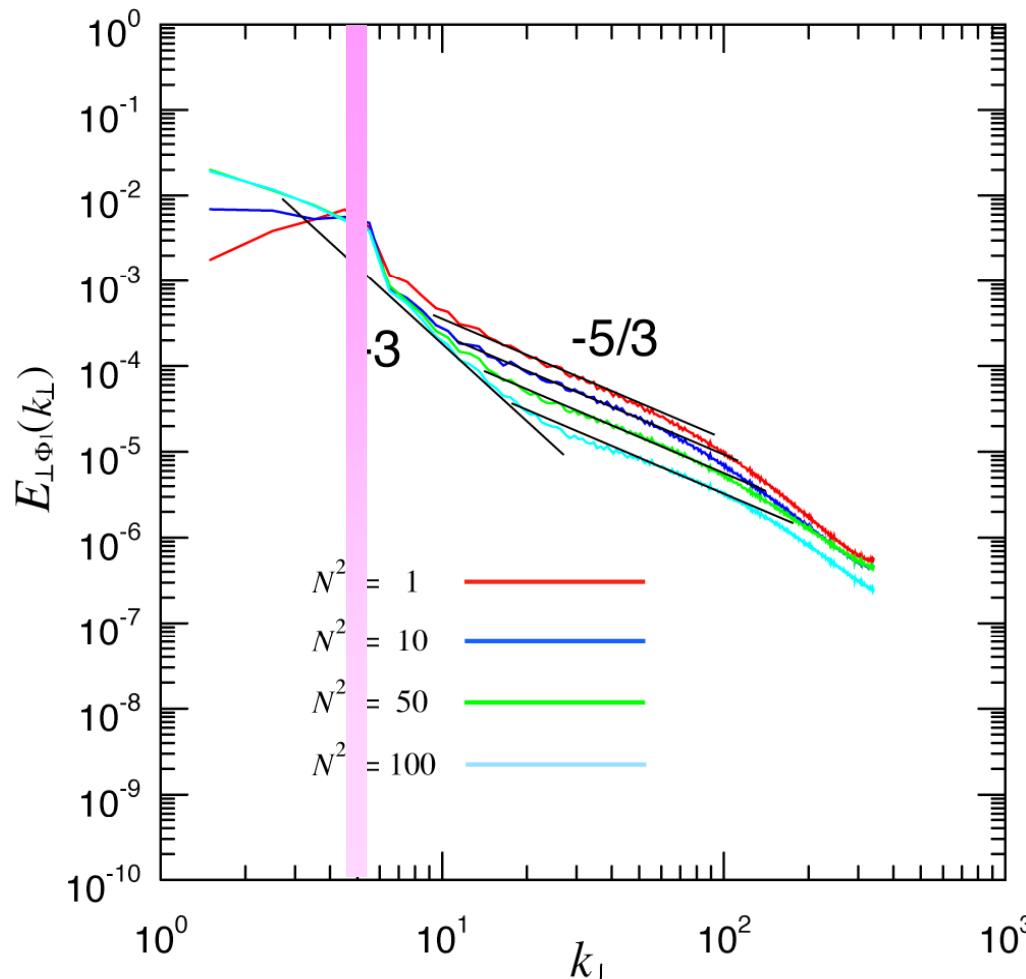
(horizontal enstrophy dissipation)

$$\varepsilon_{\perp\Phi_1} = 2\nu \int_0^\infty k_\perp^2 E_{\Phi_1}(k_\perp) dk_\perp$$

(horizontal energy dissipation)

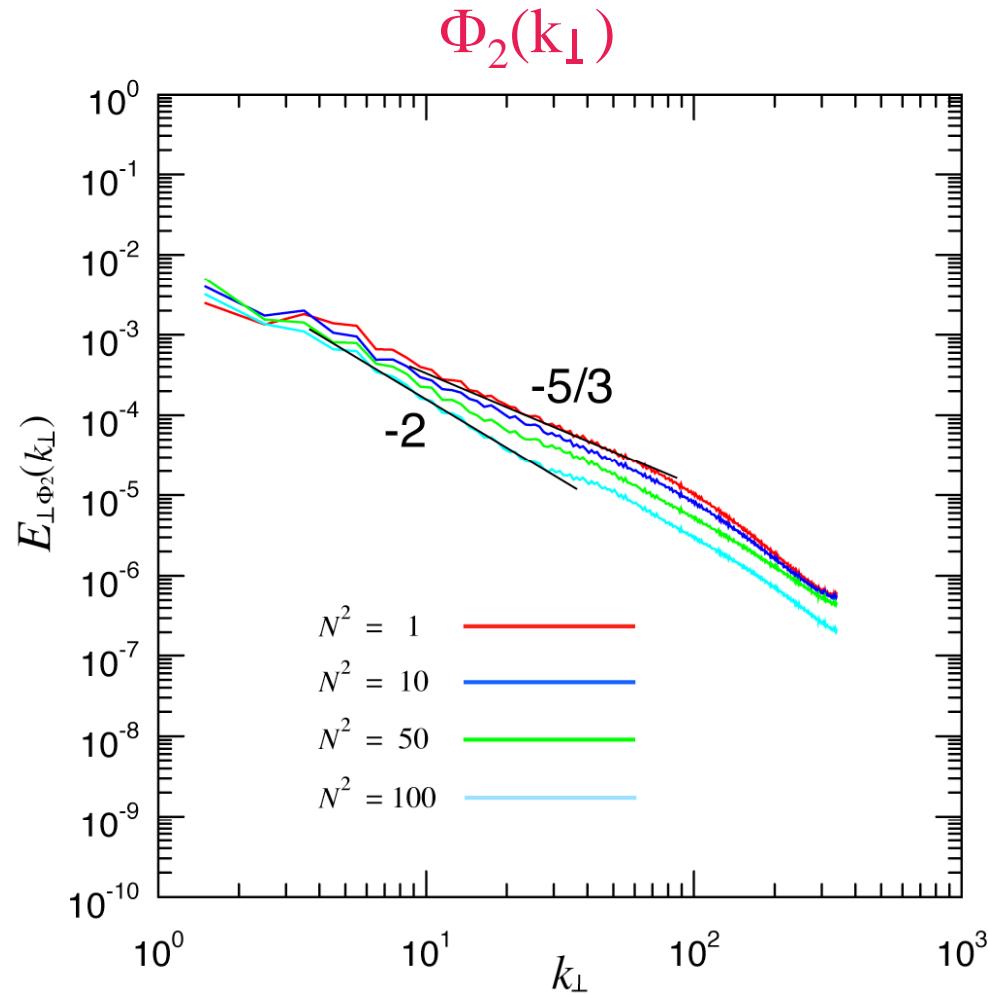
Transition wavenumber

$$\alpha \eta_{\perp\Phi_1}^{2/3} k_c^{-3} = C_K \varepsilon_{\perp\Phi_1}^{2/3} k_c^{-5/3} \quad \rightarrow \quad k_c = \left(\frac{\alpha}{C_k} \right)^{3/4} \sqrt{\frac{\eta_{\perp\Phi_1}}{\varepsilon_{\perp\Phi_1}}}$$



N^2	k_c
100	11.89
50	10.66
10	8.54
1	7.56

$\Phi_2(\mathbf{k})$ spectra for various N



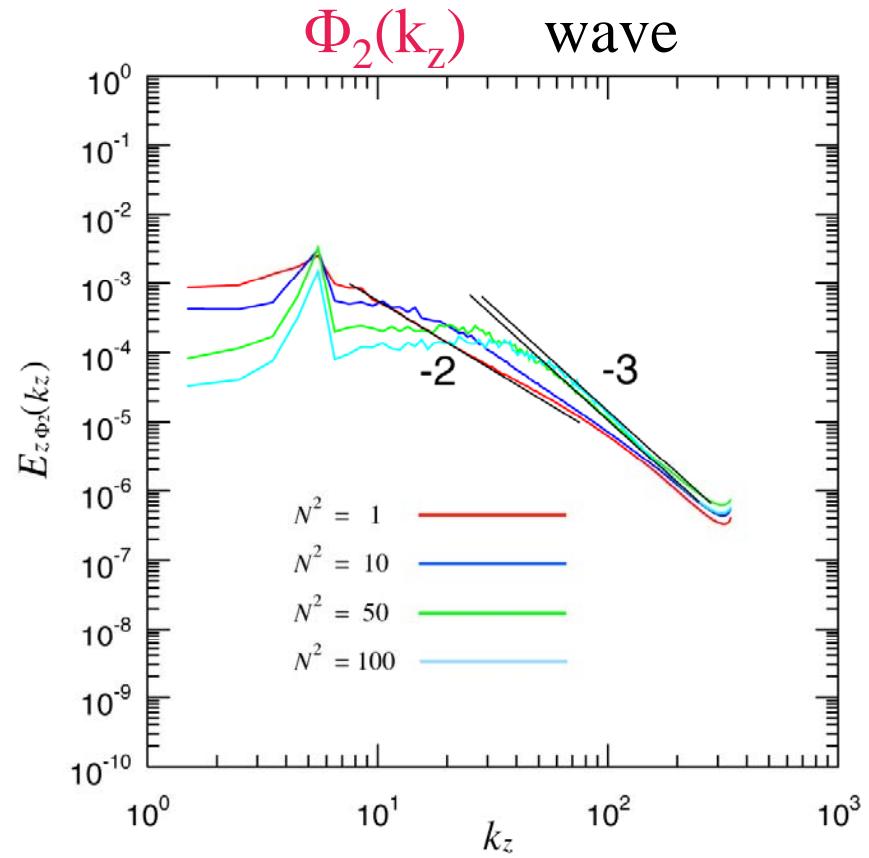
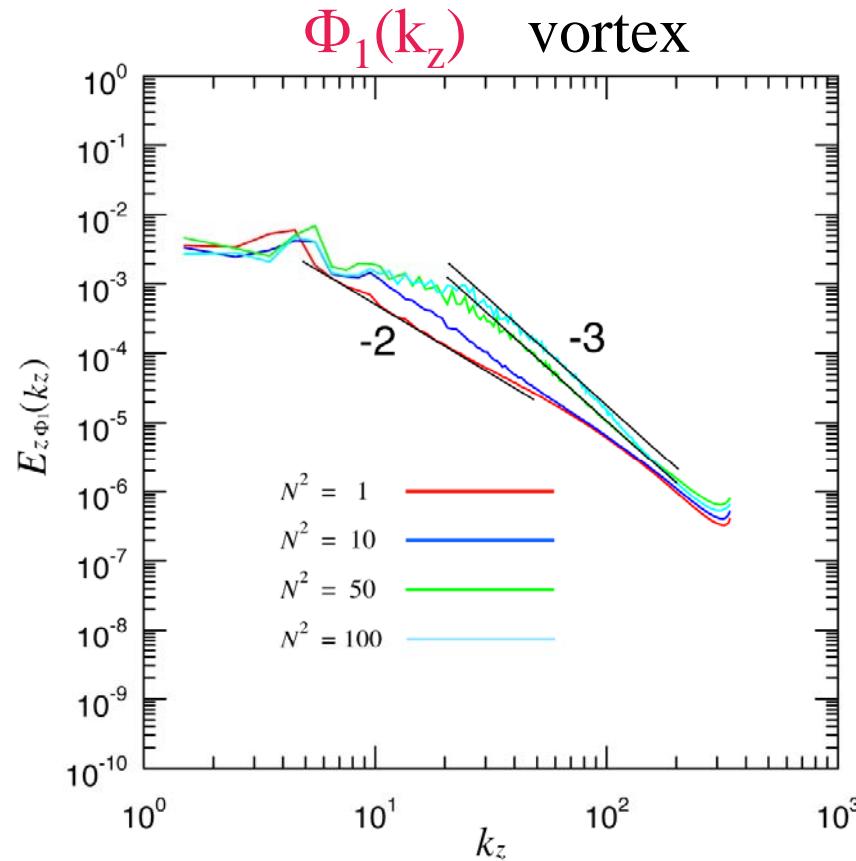
$$E_{\Phi_2}(k_\perp) = \begin{cases} \beta \sqrt{N \varepsilon_{\perp\Phi_2}} k_\perp^{-2} & (k_\perp < k_c) \\ C_K \varepsilon_{\perp\Phi_2}^{2/3} k_\perp^{-5/3} & (k_\perp > k_c) \end{cases}$$

$$k_c = \left(\frac{\beta}{C_k} \right)^3 \sqrt{\frac{N^3}{\varepsilon_{\perp\Phi_2}}}$$

transition wavenumber

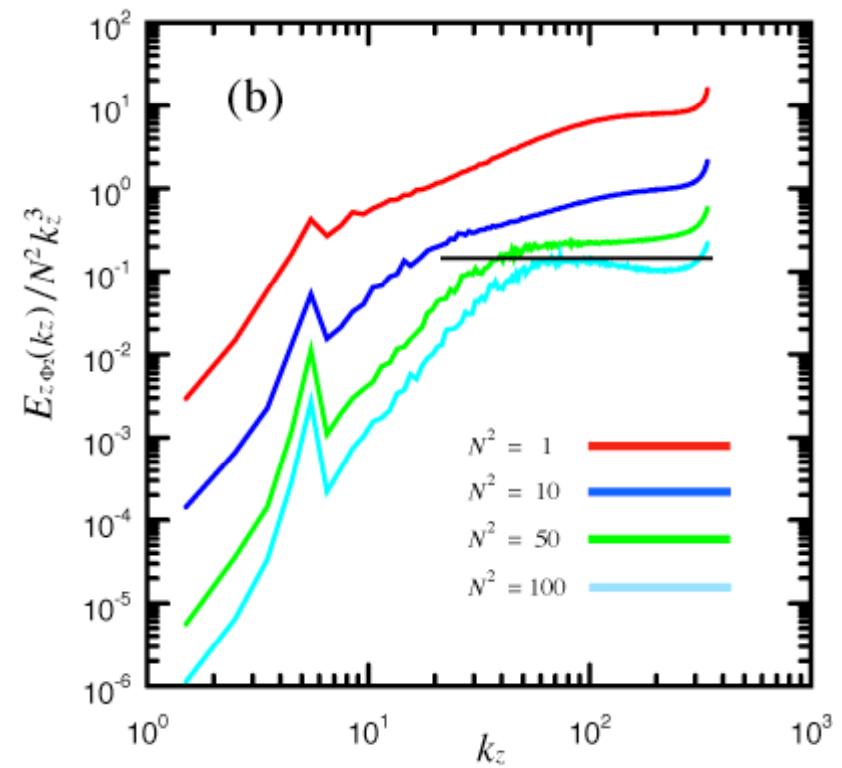
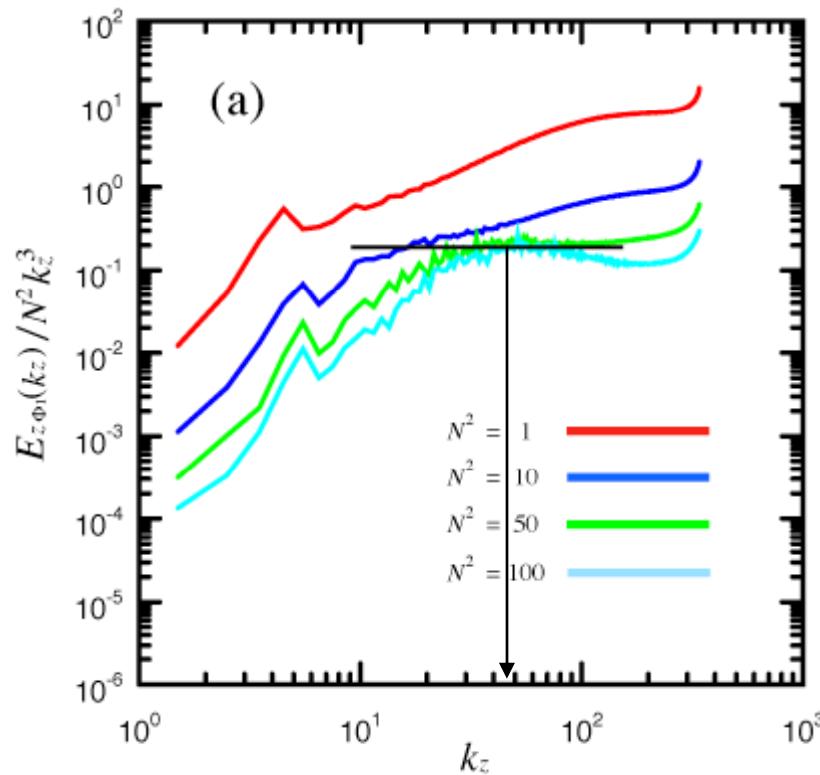
	N^2	k_c
Proportional to the Ozmidov scale (based on horizontal dissipation)	100	6.02
	50	2.63
	10	0.35
	1	3.31×10^{-2}

$\Phi_1(k_z)$ & $\Phi_2(k_z)$ spectra for various N



- ◆ flat spectra at large scales
- ◆ steep (-3) spectra at small scales for strong stratification

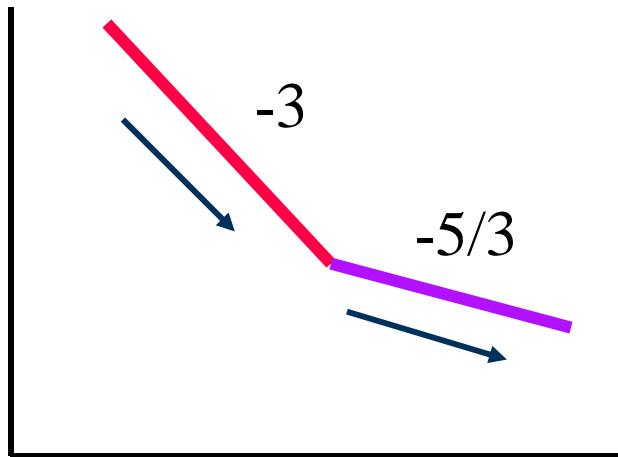
Compensated spectra for $\Phi_1(k_z)$ & $\Phi_2(k_z)$



- ◆ starting wavenumber for -3: $\sim N/u_{rms}$
- ◆ coefficient : $\sim 0.1 \rightarrow 0.1N^2 k_z^{-3}$
(saturation spectra ?)

N^2	N/u_{rms}
100	48.47
50	33.80
10	17.63
1	5.74

More than 1 inertial range



$$\frac{\partial}{\partial t} E_K(k_{\perp}) = T_K(k_{\perp}) + B(k_{\perp}) + D_K(k_{\perp}) + F_K(k_{\perp})$$

$$\left\{ \begin{array}{l} T_K(k_{\perp}) = -\text{Im} \sum_{\mathbf{k}+\mathbf{p}+\mathbf{q}=\mathbf{0}} P_{ijm}(\mathbf{k}) \tilde{u}_i(\mathbf{k}) \tilde{u}_j(\mathbf{p}) \tilde{u}_m(\mathbf{q}) \\ B(k_{\perp}) = \text{Re} \sum_{\mathbf{k}+\mathbf{p}=0} P_{i3}(\mathbf{k}) \tilde{u}_i(\mathbf{k}) \tilde{\theta}(\mathbf{p}) \\ D_K(k_{\perp}) = -\nu k^2 E_K(k_{\perp}) \end{array} \right.$$

Kolmogorov 乱流 \rightarrow

$$\int_0^{k_{\perp}} T_K(k_{\perp}) dk_{\perp} \equiv -\phi \quad (\text{flux})$$

地球流体乱流 \rightarrow

$$\int_{a_i}^{b_i} T_K(k_{\perp}) dk_{\perp} \equiv -\phi_i \quad (i^{\text{th}} \text{ flux})$$

?

Summary

- ◆ Energy spectra are investigated for stably stratified turbulence with 1024^3 pseudospectral DNS simulations.
- ◆ Horizontal spectra show clear transition from 2D to 3D Kolmogorov spectra.
- ◆ Horizontal spectra are scaled by anisotropic dissipation of energy and enstrophy.
- ◆ Vertical spectra show a flat part at large scales and tend to have steeper spectrum(-3) as N becomes large.